

ERGODICITY FOR A CLASS OF MARKOV PROCESSES AND APPLICATIONS TO RANDOMLY FORCED PDE'S. II

ARMEN SHIRIKYAN

Laboratoire de Mathématiques
Université de Paris-Sud XI, Bâtiment 425
91405 Orsay Cedex, France

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ABSTRACT. The paper is devoted to studying the problem of ergodicity for the complex Ginzburg–Landau (CGL) equation perturbed by an external random force. We show that the conditions of a simple general result established in [22] are fulfilled for the equation in question. As a consequence, we prove that the corresponding family of Markov processes has a unique stationary distribution, which possesses a mixing property. The result of this paper was announced in the joint work with Sergei Kuksin [14].

1. Introduction. The objective of this paper is to prove the uniqueness of stationary measure for the complex Ginzburg–Landau (CGL) equation perturbed by a random force. More precisely, we study the equation

$$\dot{u} - (\nu + i\alpha)\Delta u + i\beta|u|^{2\sigma}u = h(x) + \eta(t, x), \quad x \in D, \quad (1)$$

where $D \subset \mathbb{R}^n$ is a bounded domain, $h(x)$ is a deterministic function, and $\eta(t, x)$ is a random process white in time and smooth in the space variables. (See Section 2.1 for the precise assumptions imposed on the right-hand side.) Equation (1) is supplemented with the Dirichlet boundary condition and an initial condition at the time $t = 0$. We show that if the distribution of η is sufficiently non-degenerate, then the random dynamical system associated with (1) has a unique stationary measure μ , and any other solution converges to μ in distribution as $t \rightarrow \infty$.

The problem of ergodicity was studied by many authors for various classes of randomly forced PDE's. We refer the reader to the reviews [1, 10, 21] and to the Introduction of [22] for a concise summary of the results obtained and the techniques developed. Here we mention only three papers that are directly related to the equation considered in the present article. Namely, Hairer [6] studied a real Ginzburg–Landau equation on a multidimensional torus and proved the uniqueness of stationary measure and an exponential mixing property for it, Odasso [19] established similar results for a class of CGL equations with strong nonlinear dissipation, and Debussche and Odasso [2] proved uniqueness and polynomial mixing for a damped one-dimensional Schrödinger equation.

The method used in this paper is based on studying a pair of independent copies of the Markov process generated by (1). This approach was applied in [22] to give a simple proof of the uniqueness and a mixing property of stationary measure for the

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2D Navier–Stokes (NS) system in a bounded domain. The case of the CGL equation is technically more complicated. However, the main ideas remain the same, and we refer the reader to the Introduction of [22] for their informal explanation. Here we only clarify the difference between the cases of NS and CGL equations.

We wish to show that the distributions of solutions for (1) converge to a unique stationary measure in the Kantorovich–Wasserstein (KW) metric over the Sobolev space H_0^1 (see (2)). A crucial point of the proof is to estimate the distance between the distributions of solutions to (1) issued from different initial data. This is done in two steps. We first show that if the space of probability measures is endowed with the KW metric over the space L^2 (see (3)), then the arguments of [22] combined with a new a priori estimate established in Proposition 2 yield a uniform (in time) bound for the distance between solutions. The above argument does not apply to the KW metric over H_0^1 , because the a priori estimates available for higher Sobolev norms of solutions for the CGL equation are not strong enough. To overcome this difficulty, we prove that the Markov semigroup defined by (1) in the space of measures possesses a regularising property (see Proposition 4). Combining it with the bound for the KW distance over L^2 , we obtain the desired result.

Here is the plan of this paper. In Section 2, we state a well-known result on correctness of the initial-boundary value problem for Eq. (1) and establish some a priori estimates for solutions. The main result is presented in Section 3. To prove it, we show that the conditions of Theorem 1.2 in [22] are satisfied for the model in question. Finally, in the Appendix, we have compiled some auxiliary results used in the main text.

Notation. Let X be a separable Banach space with a norm $\|\cdot\|_X$. We denote by $B_X(r)$ the closed ball in X of radius r centred at the origin. We always assume that X is endowed with the Borel σ -algebra $\mathcal{B}(X)$ and denote by $\mathcal{P}(X)$ the set of probability measures on $(X, \mathcal{B}(X))$. We write $C_b(X)$ and $\mathcal{L}(X)$ for the spaces of bounded continuous and bounded Lipschitz continuous functions $f : X \rightarrow \mathbb{R}$ and endow them with the natural norms

$$\|f\|_\infty = \sup_{u \in X} |f(u)|, \quad \|f\|_{\mathcal{L}} = \sup_{u \in X} |f(u)| + \sup_{u \neq v} \frac{|f(u) - f(v)|}{\|u - v\|_X}.$$

If ξ is a random variable on a probability space $(\Omega, \mathcal{F}, \mathbb{P})$, then we denote by $\mathcal{D}(\xi)$ or $\mathbb{P}\{\xi \in \cdot\}$ the distribution of ξ . If $a, b \in \mathbb{R}$, then $a \vee b$ ($a \wedge b$) stands for the maximum (minimum) of a and b .

We deal with the spaces H , H^1 , and H^2 introduced in Subsection 2.1 together with corresponding norms $\|\cdot\|$, $\|\cdot\|_1$, and $\|\cdot\|_2$. If $f \in C_b(H^1)$ and $\mu \in \mathcal{P}(H^1)$, then (f, μ) denotes the integral of f over H^1 with respect to μ . If $\mu_1, \mu_2 \in \mathcal{P}(H^1)$, then we write*

$$\|\mu_1 - \mu_2\|_{\mathcal{L}}^* := \sup_{\|f\|_{\mathcal{L}} \leq 1} |(f, \mu_1) - (f, \mu_2)|, \quad (2)$$

$$|\mu_1 - \mu_2|_{\mathcal{L}}^* := \sup_{\|g\|_{\mathcal{L}} \leq 1} |(g, \mu_1) - (g, \mu_2)|, \quad (3)$$

where $f \in \mathcal{L}(H^1)$, $g \in \mathcal{L}(H)$, and we denote by $\|\cdot\|_{\mathcal{L}}$ and $|\cdot|_{\mathcal{L}}$ the norms in the spaces $\mathcal{L}(H^1)$ and $\mathcal{L}(H)$, respectively.

*Any measure $\mu \in \mathcal{P}(H^1)$ can be extended to a measure on H by the formula $\mu(\Gamma) = \mu(\Gamma \cap H^1)$, and therefore the integral (g, μ) is well defined for any $g \in C_b(H)$.

2. Initial-boundary value problem.

2.1. Well-posedness and a priori estimates. Let $D \subset \mathbb{R}^n$, $n \geq 1$, be a bounded domain with smooth boundary ∂D . Consider the problem

$$\dot{u} - (\nu + i\alpha)\Delta u + i\beta|u|^{2\sigma}u = h(x) + \eta(t, x), \quad (4)$$

$$u|_{\partial D} = 0, \quad (5)$$

$$u(0, x) = u_0(x). \quad (6)$$

Here $\nu > 0$, $\alpha > 0$, $\beta \geq 0$, and $\sigma \geq 0$ are some constants, $h(x)$ is a deterministic function belonging to the Sobolev space of order 1, and $\eta(t, x)$ is a random process white in time and smooth in x . More precisely, we assume that

$$\eta(t, x) = \frac{\partial}{\partial t}\zeta(t, x), \quad \zeta(t, x) = \sum_{j=1}^{\infty} b_j \beta_j(x) e_j(x), \quad (7)$$

where $\{\beta_j = \beta_j^1 + i\beta_j^2\}$ is a sequence of complex-valued independent Brownian motions defined on a probability space $(\Omega, \mathcal{F}, \mathbb{P})$ with right-continuous filtration \mathcal{F}_t , $\{e_j\}$ is a complete set of eigenvectors for the Dirichlet Laplacian in D with eigenvalues $\alpha_1 < \alpha_2 \leq \alpha_3 \leq \dots$, and $b_j \geq 0$ are real constants such that

$$B_0 := \sum_{j=1}^{\infty} b_j^2 < \infty. \quad (8)$$

Let H be the space of complex-valued square-integrable functions on D . We shall regard it as a real Hilbert space with the scalar product

$$(u, v) = \operatorname{Re} \int_D u(x) \overline{v(x)} dx,$$

and the corresponding norm will be denoted by $\|u\|$. Let $H^1 = H_0^1(D, \mathbb{C})$ and $H^2 = H^2(D, \mathbb{C}) \cap H^1$, where $H^s(D, \mathbb{C})$ stands for the Sobolev space of order s in the domain D and $H_0^1(D, \mathbb{C})$ for the space of functions $u \in H^1(D, \mathbb{C})$ vanishing on ∂D . The spaces H^1 and H^2 are endowed with the norms

$$\|u\|_1 = \left(\int_D |\nabla u(x)|^2 dx \right)^{\frac{1}{2}}, \quad \|u\|_2 = \left(\int_D |\Delta u(x)|^2 dx \right)^{\frac{1}{2}}.$$

From now on, we assume that

$$0 \leq \sigma \leq \frac{2}{n-2} \quad \text{for } n \geq 3, \quad \sigma \geq 0 \quad \text{for } n = 1, 2. \quad (9)$$

Let us define the following continuous functionals on H^1 :

$$\mathcal{H}_0(u) := \frac{1}{2} \|u\|^2 = \frac{1}{2} \int_D |u(x)|^2 dx,$$

$$\mathcal{H}_1(u) := \int_D \left(\frac{\alpha}{2} |\nabla u(x)|^2 + \frac{\beta}{2\sigma+2} |u(x)|^{2\sigma+2} \right) dx.$$

If X is Banach space and $J \subset \mathbb{R}$ is a closed interval, then we denote by $C(J, X)$ the space of continuous functions $f : J \rightarrow X$ and by $L_{\text{loc}}^2(J, X)$ the space of measurable functions $f : J \rightarrow X$ such that $\int_I \|f(t)\|_X^2 dt < \infty$ for any finite subinterval $I \subset J$.

The following theorem establishes the well-posedness of (4) – (7). Its proof is carried out by standard methods and can be found in [14] for the case $h \equiv 0$, $\sigma = 1$, and $1 \leq n \leq 4$. We refer the reader to [9, 18] for more general existence and uniqueness results for SPDE's.

Theorem 1. *Let u_0 be an H^1 -valued \mathcal{F}_0 -measurable random variable such that $\mathbb{E} \mathcal{H}_1(u_0) < \infty$. Suppose that $h \in H^1(D, \mathbb{C})$, inequalities (9) are satisfied, and*

$$B_1 := \sum_{j=1}^{\infty} \alpha_j b_j^2 < \infty, \quad M := \sum_{j=1}^{\infty} b_j^2 \|e_j\|_{L^\infty}^2 < \infty. \quad (10)$$

Then the following statements hold.

- (i) *There is an \mathcal{F}_t -adapted random process $u(t) = u(t, x)$, $t \geq 0$, whose almost every trajectory belongs to the space*

$$\mathcal{X} := C(\mathbb{R}_+; H^1) \cap L_{\text{loc}}^2(\mathbb{R}_+; H^2)$$

and satisfies Eqs. (4) and (6) in the sense that

$$u(t) = u_0 + \int_0^t ((\nu + i\alpha)\Delta u(s) - i\beta|u(s)|^{2\sigma}u(s)) ds + th + \zeta(t), \quad t \geq 0,$$

where the left- and right-hand sides are regarded as elements of H .

- (ii) *The process $u(t)$ constructed in (i) is unique in the sense that if $\tilde{u}(t)$ is another random process satisfying (i), then, with probability 1, we have $u(t) = \tilde{u}(t)$ for all $t \geq 0$.*
- (iii) *The random process $\mathcal{H}_0(u(t))$ and $\mathcal{H}_1(u(t))$ possess stochastic differentials, which have the form*

$$d\mathcal{H}_0(u(t)) = (-\nu \|u(t)\|_1^2 + B_0 + (u(t), h)) dt + (u(t), d\zeta(t)), \quad (11)$$

$$\begin{aligned} d\mathcal{H}_1(u(t)) = & \left(\alpha B_1 + \beta(\sigma + 1) \sum_{j=1}^{\infty} b_j^2 (|u|^{2\sigma}, e_j^2) \right. \\ & - \nu \{ \alpha \|u\|_2^2 + \beta(\sigma + 1)(|u|^{2\sigma}, |\nabla u|^2) + \beta\sigma(|u|^{2(\sigma-1)}u^2, (\nabla u)^2) \} \\ & \left. + (-\alpha\Delta u + \beta|u|^{2\sigma}u, h) \right) dt + (-\alpha\Delta u + \beta|u|^{2\sigma}u, d\zeta), \end{aligned} \quad (12)$$

where the constant B_0 is defined in (8). Moreover, for any $t \geq 0$, we have

$$\mathbb{E} \|u(t)\|^2 + \nu \int_0^t \mathbb{E} \|u(s)\|_1^2 ds \leq \mathbb{E} \|u_0\|^2 + (B_0 + (\alpha_1\nu)^{-1}\|h\|^2) t, \quad (13)$$

$$\begin{aligned} \mathbb{E} \mathcal{H}_1(u(t)) + \frac{\nu}{2} \int_0^t \mathbb{E} \{ \alpha \|u(s)\|_2^2 + \beta(|u(s)|^{2\sigma}, |\nabla u|^2) \} ds \\ \leq \mathbb{E} \mathcal{H}_1(u_0) + C(B_1 + \nu^{-(2\sigma+1)}\|h\|_1^{2\sigma+2} + 1) t, \end{aligned} \quad (14)$$

where $C > 0$ is a constant depending only on α and β .

2.2. Exponential martingale inequalities. For any function $u(t)$ belonging to the space \mathcal{X} (see Theorem 1), we set

$$\mathcal{E}_u(t) = \|u(t)\|^2 + \nu \int_0^t \|u(s)\|_1^2 ds.$$

Proposition 1. *Suppose that (9) and (10) holds. Then there exist positive constants K and γ such that the solution $u(t)$ of (4) – (6) constructed in Theorem 1 satisfies the inequality*

$$\mathbb{P} \left\{ \sup_{t \geq 0} (\mathcal{E}_u(t) - Kt) \geq \|u_0\|^2 + \rho \right\} \leq e^{-\gamma\rho} \quad \text{for any } \rho > 0. \quad (15)$$

Proof. We only outline the proof, which repeats the arguments used in [16, 13]. In view of (11), we have

$$\mathcal{H}_0(u(t)) + \nu \int_0^t \|u(s)\|_1^2 ds = \mathcal{H}_0(u_0) + B_0 t + \int_0^t (u(s), h) ds + \int_0^t (u(s), d\zeta(s)).$$

Setting $K = 2(B_0 + (\nu\alpha_1)^{-1}\|h\|^2)$, we easily show that

$$\mathcal{E}_u(t) \leq \|u_0\|^2 + Kt + 2 \int_0^t (u(s), d\zeta(s)) - \frac{\nu}{2} \int_0^t \|u(s)\|_1^2 ds. \quad (16)$$

Let us define a martingale by the formula

$$\mathcal{M}_t = 2 \int_0^t (u(s), d\zeta(s))$$

and note that its quadratic variation satisfies the inequality

$$\langle \mathcal{M} \rangle_t = 4 \sum_{j=1}^{\infty} b_j^2 \int_0^t \{(u(s), e_j)^2 + (u(s), ie_j)^2\} ds \leq \frac{4b^2}{\alpha_1} \int_0^t \|u(s)\|_1^2 ds,$$

where $b = \max_j b_j$. Combining this inequality with (16), we obtain

$$\mathcal{E}_u(t) \leq \|u_0\|^2 + Kt + \left(\mathcal{M}_t - \frac{\gamma \langle \mathcal{M} \rangle_t}{2} \right),$$

where $\gamma = \frac{\alpha_1 \nu}{4b^2}$. It follows that

$$\begin{aligned} \mathbb{P} \left\{ \sup_{t \geq 0} (\mathcal{E}_u(t) - Kt) - \|u_0\|^2 \geq \rho \right\} &\leq \mathbb{P} \left\{ \sup_{t \geq 0} \left(\mathcal{M}_t - \frac{\gamma \langle \mathcal{M} \rangle_t}{2} \right) \geq \rho \right\} \\ &\leq \mathbb{P} \left\{ \sup_{t \geq 0} \exp(\gamma \mathcal{M}_t - \frac{\gamma^2 \langle \mathcal{M} \rangle_t}{2}) \geq \rho \right\}. \end{aligned} \quad (17)$$

Since $\exp(\gamma \mathcal{M}_t - \gamma^2 \langle \mathcal{M} \rangle_t / 2)$ is a supermartingale with mean value not exceeding 1, the classical supermartingale inequality (see Theorem III.6.11 in [17]) implies that the probability on the right-hand side of (17) can be estimated from above by $e^{-\gamma\rho}$. \square

We now consider the functional

$$\mathcal{J}_u(t) = \mathcal{H}_1(u(t)) + \nu \mathcal{I}_u(t),$$

where

$$\mathcal{I}_u(t) = \frac{1}{2} \int_0^t \{ \alpha \|u(s)\|_2^2 + \beta (|u(s)|^{2\sigma}, |\nabla u(s)|^2) \} ds.$$

Proposition 2. *Suppose that (10) holds and that $\sigma \leq \frac{2}{n-2} \wedge 1$. Then there is a constant $p \geq 2$ such that for any $T > 0$ and $\rho > 0$ the solution $u(t)$ of (4) – (6) satisfies the inequality*

$$\mathbb{P} \left\{ \sup_{0 \leq t \leq T} \mathcal{J}_u(t) \geq \mathcal{H}_1(u_0) + C \|u_0\|^p + K + \rho \right\} \leq \exp(-\gamma\rho^{2/p}), \quad (18)$$

where C , K and γ are positive constants not depending on u_0 and ρ .

Proof. We repeat the scheme used in the proof of Proposition 1. The difference is that the quadratic variation of the corresponding martingale cannot be estimated by $\mathcal{I}_u(t)$.

Step 1. In view of (12) and the inequality

$$|(|u|^{2(\sigma-1)}u^2, (\nabla u)^2)| \leq (|u|^{2\sigma}, |\nabla u|^2),$$

we have

$$\mathcal{H}_1(u(t)) + 2\nu\mathcal{I}_u(t) \leq \mathcal{H}_1(u_0) + F_t + \mathcal{M}_t, \quad (19)$$

where

$$F_t = \alpha B_1 t + \int_0^t (-\alpha\Delta u + \beta|u|^{2\sigma}u, h) ds + \beta(\sigma + 1) \sum_{j=1}^{\infty} b_j^2 \int_0^t (|u|^{2\sigma}, e_j^2) ds,$$

$$\mathcal{M}_t = \int_0^t (-\alpha\Delta u + \beta|u|^{2\sigma}u, d\zeta).$$

Suppose we have found $p \geq 2$ such that

$$|F_t| \leq \frac{\nu}{2} \mathcal{I}_u(t) + C_1 t, \quad (20)$$

$$\langle \mathcal{M} \rangle_t \leq C_2 \mathcal{I}_u(t) + C_3 \int_0^t \|u\|^p ds. \quad (21)$$

Here and henceforth, we denote by C_i unessential positive constants. Combining (19) – (21), we see that

$$\sup_{0 \leq t \leq T} \mathcal{J}_u(t) \leq \mathcal{H}_1(u_0) + C_4 + C_5 \sup_{0 \leq t \leq T} \|u(t)\|^p + \sup_{0 \leq t \leq T} \left(\mathcal{M}_t - \frac{c\langle \mathcal{M} \rangle_t}{2} \right), \quad (22)$$

where $c > 0$ is sufficiently small. In view of (15), there are positive constants K_1 and γ_1 such that

$$\mathbb{P} \left\{ \sup_{0 \leq t \leq T} \|u(t)\|^p \geq \|u_0\|^p + K_1 + \rho \right\} \leq \exp(-\gamma_1 \rho^{2/p}). \quad (23)$$

Furthermore, as in the proof of Proposition 1, the supermartingale inequality implies that

$$\mathbb{P} \left\{ \sup_{0 \leq t \leq T} \left(\mathcal{M}_t - \frac{c\langle \mathcal{M} \rangle_t}{2} \right) \right\} \leq \exp(-c\rho). \quad (24)$$

Comparing (22) – (24), we arrive at the desired inequality (18).

Step 2. We now prove (20) and (21). Using the Hölder inequality, the Sobolev embedding theorems, and the second condition in (10), we derive

$$\begin{aligned} \left| \int_0^t (-\alpha\Delta u + \beta|u|^{2\sigma}u, h) ds \right| &\leq \int_0^t \left(\alpha\|u\|_2 \|h\| + \beta\|u\|_{L^{2\sigma+2}}^{2\sigma+1} \|h\|_{L^{2\sigma+2}} \right) ds \\ &\leq \frac{\nu}{4} \mathcal{I}_u(t) + C_6 (\nu^{-1} \|h\|^2 + \nu^{-(2\sigma+1)} \|h\|_1^{2\sigma+2}) t, \\ \sum_{j=1}^{\infty} b_j^2 \int_0^t (|u|^{2\sigma}, e_j^2) ds &\leq M \int_0^t \|u\|_{L^{2\sigma}}^{2\sigma} ds \leq \frac{\nu}{4\beta(\sigma+1)} \mathcal{I}_u(t) + C_7 t. \end{aligned}$$

These inequalities imply (20).

To prove (21), we write

$$\begin{aligned} \langle \mathcal{M} \rangle_t &= \sum_{j=1}^{\infty} b_j^2 \int_0^t \left\{ (-\alpha\Delta u + \beta|u|^{2\sigma}u, e_j)^2 + (-\alpha\Delta u + \beta|u|^{2\sigma}u, ie_j)^2 \right\} ds \\ &\leq 2\alpha^2 B_0 \int_0^t \|u\|_2^2 ds + 4\beta^2 M \int_0^t \|u\|_{L^{2\sigma+1}}^{4\sigma+2} ds. \end{aligned} \quad (25)$$

Thus, the required inequality (21) will be established if we show that

$$\|u\|_{L^{2\sigma+1}}^{4\sigma+2} \leq (|u|^{2\sigma}, |\nabla u|^2) + C_8 \|u\|^p. \quad (26)$$

To simplify formulas, we confine ourselves to the case $n = 4$ and $\sigma = 1$. In view of the Gagliardo–Nirenberg inequality (see Theorem 6.4.1 in [7]), we have

$$\|u\|_{L^3}^6 = \|u^2\|_{L^{3/2}}^3 \leq C_9 \|u^2\|_{L^1}^{5/3} \|u^2\|_1^{4/3} \leq C_{10} \|u\|^{10/3} (|u|^2, |\nabla u|^2)^{2/3}.$$

This implies inequality (26) with $p = 10$. \square

3. Uniqueness of stationary distribution and mixing.

3.1. Main result. Throughout this section, we shall assume that the parameter $\sigma \geq 0$ satisfies the inequalities (cf. (9))

$$\sigma \leq \frac{2}{n} \quad \text{for } n \geq 3, \quad \sigma < 1 \quad \text{for } n = 2, \quad \sigma \leq 1 \quad \text{for } n = 1. \quad (27)$$

Let us denote by (u_t, \mathbb{P}_u) the family of Markov processes associated with the problem (4) – (6) and parametrised by the deterministic initial condition $u_0 = u \in H^1$. Let

$$P(t, u, \Gamma) = \mathbb{P}_u\{u_t \in \Gamma\}, \quad u \in H^1, \quad \Gamma \in \mathcal{B}(H^1),$$

be the transition function for the family (u_t, \mathbb{P}_u) and let $P_t : C_b(H^1) \rightarrow C_b(H^1)$ and $P_t^* : \mathcal{P}(H^1) \rightarrow \mathcal{P}(H^1)$ be the corresponding Markov semigroups. Recall that $\mu \in \mathcal{P}(H^1)$ is called a *stationary measure* for (u_t, \mathbb{P}_u) if $P_t^* \mu = \mu$ for all $t \geq 0$. The following theorem is the main result of this paper.

Theorem 2. *Suppose that $h \in H^1(D, \mathbb{C})$, conditions (10) and (27) are satisfied, and*

$$b_j \neq 0 \quad \text{for all } j \geq 1. \quad (28)$$

Then for any $\nu > 0$ the Markov family (u_t, \mathbb{P}_u) has a unique stationary measure. Moreover, the measure μ is mixing in the sense that, for any $\lambda \in \mathcal{P}(H^1)$, we have

$$\|P_t^* \lambda - \mu\|_{\mathcal{L}}^* \rightarrow 0 \quad \text{as } t \rightarrow \infty.$$

A proof of Theorem 2 is given in the next subsection. Here we outline the main ideas. To simplify the presentation, in what follows we confine ourselves to the case $n = 3$ or 4 , which is the most difficult.

Let $(\mathbf{u}_t, \mathbb{P}_{\mathbf{u}})$ be a pair of independent copies of the family (u_t, \mathbb{P}_u) . In other words, $(\mathbf{u}_t, \mathbb{P}_{\mathbf{u}})$ is a family of Markov processes in $\mathbf{H}^1 = H^1 \times H^1$ whose transition function is given by the formula

$$\mathbf{P}(t, \mathbf{u}, \Gamma \times \Gamma') = P(t, u, \Gamma) P(t, u', \Gamma')$$

for any $\mathbf{u} = (u, u') \in \mathbf{H}$ and $\Gamma, \Gamma' \in \mathcal{P}(H^1)$. Let $\mathbf{G}_m = B_{H^1}(1/m) \times B_{H^1}(1/m)$, where $B_{H^1}(r)$ stands for the closed ball in H^1 of radius r centred at the origin. Denote by τ_m the first hitting time of \mathbf{G}_m :

$$\tau_m = \min\{t \geq 0 : \mathbf{u}_t \in \mathbf{G}_m\}.$$

In view of Theorem 1.2 in [22], the required result will be established if we show that the following two properties hold.

(P₁) For any $\mathbf{u} = (u, u') \in \mathbf{H}^1$ and $m \geq 1$, we have

$$\mathbb{P}_{\mathbf{u}}\{\tau_m < \infty\} = 1.$$

(P₂) There is a constant $T > 0$ and sequence $\delta_m > 0$ going to zero as $m \rightarrow \infty$ such that

$$\sup_{t \geq T} \|P(t, u, \cdot) - P(t, u', \cdot)\|_{\mathcal{L}}^* \leq \delta_m \quad \text{for any } \mathbf{u} \in \mathbf{G}_m. \quad (29)$$

Proof of (P₁) literally repeats the arguments used in [22] for the case of the 2D Navier–Stokes system, and we shall not dwell on it. Let us sketch the proof of (P₂) (cf. Step 4 in [22, Section 3.1]).

Let us denote by $(\Omega, \mathcal{F}, \mathbb{P})$ the probability space associated with the problem (4) – (7), and let $u_t(\omega, u)$ be the solution of (4), (5) issued from the initial point $u \in H^1$. The proof of (P₂) is based on the two propositions below. The first of them enables one to estimate the distance between the distributions of solutions for (4), (5) with different initial points. This type of results were obtained in [11, 12] for discrete-time forces and in [4] for white noise and were developed later in a number of works. Our presentation is close to that of the papers [15, 6], in which the closeness of distributions is described with the help of transformation of the underlying probability space.

Proposition 3. *For any $\delta > 0$ there is $\varepsilon > 0$ such that, for any $u \in B_{H^1}(\varepsilon)$, one can find a measurable transformation $\Psi_u : \Omega \rightarrow \Omega$ satisfying the inequalities*

$$\sup_{t \geq 0} \mathbb{P} \{ \|u_t(\omega, u) - u_t(\Psi_u(\omega), 0)\| \geq \delta \} \leq \delta, \quad (30)$$

$$\|\mathbb{P} - \Psi_{u*}(\mathbb{P})\|_{\text{var}} \leq \delta, \quad (31)$$

where $\Psi_{u*}(\mathbb{P})$ stands for the image of \mathbb{P} under the transformation Ψ_u and $\|\cdot\|_{\text{var}}$ is the total variation norm.

In other words, if the initial function $u \in H^1$ is sufficiently small, then with high probability the solution is close (in the L^2 -norm) to the trajectory starting from zero and corresponding to a different value of the random parameter, which is denoted by $\Psi_u(\omega)$. Moreover, the transformation $\omega \mapsto \Psi_u(\omega)$ almost preserves the probability measure \mathbb{P} .

Using (30) and (31), one easily shows that

$$\sup_{t \geq 0} |\mathcal{D}(u_t) - \mathcal{D}(\tilde{u}_t)|_{\mathcal{L}}^* \leq 5\delta, \quad (32)$$

where $u_t = u_t(\omega, u)$, $\tilde{u}_t = u_t(\omega, 0)$, and $|\cdot|_{\mathcal{L}}^*$ stands for the dual Lipschitz metric for the norm in H (see Notation). It follows from (32) that, for any initial functions $u, u' \in B_{H^1}(\varepsilon)$, the distributions of the corresponding solutions u_t and u'_t satisfy the inequality

$$\sup_{t \geq 0} |\mathcal{D}(u_t) - \mathcal{D}(u'_t)|_{\mathcal{L}}^* \leq 10\delta, \quad (33)$$

where $\delta = \delta(\varepsilon) > 0$ goes to zero with ε .

Inequality (33) and a priori estimates for solutions imply that (29) will be established if we prove the continuity of the Markov semigroup P_t^* for appropriate norms. For any measure $\mu \in \mathcal{P}(H^1)$, we set

$$\mathcal{H}_1(\mu) = \int_{H^1} \mathcal{H}_1(u) \mu(du).$$

Proposition 4. *For any positive constants C and γ there is $\delta > 0$ such that if measures $\mu_1, \mu_2 \in \mathcal{P}(H^1)$ satisfy the inequalities*

$$\mathcal{H}_1(\mu_1) \vee \mathcal{H}_1(\mu_2) \leq C, \quad \|\mu_1 - \mu_2\|_{\mathcal{L}}^* \leq \delta, \quad (34)$$

then

$$\|P_1^* \mu_1 - P_1^* \mu_2\|_{\mathcal{L}}^* \leq \gamma. \quad (35)$$

Combining Proposition 4 with inequality (33), we conclude that Property (P₂) holds with $T = 1$.

The rest of this section is organised as follows. In Section 3.2, we give a detailed proof of the fact that Propositions 3 and 4 imply Property (P₂). Propositions 3 and 4 are established in Sections 3.3 and 3.4, respectively.

3.2. Verification of Property (P₂). *Step 1.* Let $\mathbf{u} = (u, u') \in \mathbf{G}_m$ and let u_t and u'_t be the solutions of (4), (5) issued from u and u' , respectively. To prove (29) with $T = 1$, it suffices to show that

$$\sup_{t \geq 0} |\mathcal{D}(u_t) - \mathcal{D}(u'_t)|_{\mathcal{L}}^* \rightarrow 0 \quad \text{as } m \rightarrow \infty, \quad (36)$$

and the convergence is uniform with respect to $\mathbf{u} \in \mathbf{G}_m$. Indeed, it follows from (14) and the Gronwall inequality that

$$\sup_{t \geq 0} \mathcal{H}_1(v_t) \leq C \quad \text{for } v \in B_{H^1}(1),$$

where v_t stands for the solution of (4), (5) with the initial condition v . Therefore, applying Proposition 4 to the measures $\mu_1 = \mathcal{D}(u_t)$ and $\mu_2 = \mathcal{D}(u'_t)$ with an arbitrary $t \geq 0$ and using (36), we conclude that

$$\sup_{t \geq 1} \|\mathcal{D}(u_t) - \mathcal{D}(u'_t)\|_{\mathcal{L}}^* \rightarrow 0 \quad \text{as } m \rightarrow \infty \text{ uniformly in } \mathbf{u} \in \mathbf{G}_m.$$

This is equivalent to (29).

Step 2. Convergence (36) will be established if we show that

$$\sup_{t \geq 0} |\mathcal{D}(u_t) - \mathcal{D}(\tilde{u}_t)|_{\mathcal{L}}^* \rightarrow 0 \quad \text{as } \|u\|_1 \rightarrow 0, \quad (37)$$

where \tilde{u}_t stands for the solution of (4), (5) with zero initial condition. To prove this, we use Proposition 3. Applying Lemma 3.4 in [22] to the random variables $u_t(\omega)$ and $\tilde{u}_t(\omega)$ and to the transformation Ψ_u defined in Proposition 3, we see that

$$\sup_{t \geq 0} |\mathcal{D}(u_t) - \mathcal{D}(\tilde{u}_t)|_{\mathcal{L}}^* \leq 3\delta + 2\|\mathbb{P} - \Psi_{u*}(\mathbb{P})\|_{\text{var}} \quad \text{for any } u \in B_{H^1}(\varepsilon),$$

where $\varepsilon = \varepsilon(\delta) > 0$ is sufficiently small. Combining this inequality with (31), we conclude that (32) holds for $u \in B_{H^1}(\varepsilon)$. Since $\delta > 0$ is arbitrary, we arrive at (37).

3.3. Proof of Proposition 3. We first note that the underlying probability space $(\Omega, \mathcal{F}, \mathbb{P})$ plays no role in the statement of Theorem 2. Therefore, we can assume from the very beginning that it possesses the following properties:

- Ω coincides with the space of functions $u \in C(\mathbb{R}_+, H)$ that vanish at $t = 0$;
- Ω is endowed with the topology of uniform convergence on the compact intervals $J \subset \mathbb{R}_+$, and $\mathcal{B}(\Omega)$ stands for the Borel σ -algebra on Ω ;
- \mathbb{P} is the distribution of the random process ζ defined in (7) and \mathcal{F} is the completion of $\mathcal{B}(\Omega)$ with respect to \mathbb{P} .

In this case, we can assume, without loss of generality, that

$$\zeta(t) = \omega_t \quad \text{for all } \omega \in \Omega \text{ and } t \geq 0.$$

The proof of Proposition 4 is divided into several steps.

Step 1: Construction of Ψ_u . We shall need an auxiliary result on solvability of the projection of (4) to subspaces of finite codimension. For any $N \geq 1$, denote

by H_N the $2N$ -dimensional vector space spanned by $\{e_j, ie_j, 1 \leq j \leq N\}$ and by H_N^\perp its orthogonal complement in H . Consider the problem

$$\dot{w} - (\nu + i\alpha)\Delta w + F_N^\perp(v + w) = \mathbf{Q}_N h + \mathbf{Q}_N \eta, \quad (38)$$

$$w(0) = w_0. \quad (39)$$

Here $w_0 \in H_N^\perp$ and $v \in C(\mathbb{R}_+, H_N)$ are given functions, $F_N^\perp : H^1 \rightarrow H_N^\perp$ is a continuously differentiable function defined as $F_N^\perp(u) = i\beta \mathbf{Q}_N(|u|^{2\sigma}u)$, and we denote by \mathbf{P}_N and \mathbf{Q}_N the orthogonal projections in H onto the subspaces H_N and H_N^\perp , respectively.

Lemma 1. *Under the conditions of Theorem 1, there is a set $\Omega_0 \in \mathcal{F}$ of full measure such that the following assertions hold for any $\omega \in \Omega_0$.*

- (i) *For any $v \in C(\mathbb{R}_+, H_N)$ and $w_0 \in H_N^\perp$, problem (38), (39) has a unique solution w_t that belongs to the space*

$$\mathcal{X}_N := C(\mathbb{R}_+, H^1 \cap H_N^\perp) \cap L_{\text{loc}}^2(\mathbb{R}_+, H^2 \cap H_N^\perp).$$

- (ii) *The function w_t depends only on the restrictions of v and $\mathbf{Q}_N \omega$ to the interval $[0, t]$.*

- (iii) *Let $\mathcal{W}(\cdot, \cdot, \omega) : C(\mathbb{R}_+, H_N) \times H_N^\perp \rightarrow \mathcal{X}_N$ be the operator defined as*

$$\mathcal{W}(v, w_0, \omega) = \begin{cases} w & \text{if } \omega \in \Omega_0, \\ 0 & \text{if } \omega \in \Omega \setminus \Omega_0, \end{cases}$$

where $w \in \mathcal{X}_N$ is the solution of (38), (39). Then \mathcal{W} is uniformly Lipschitz continuous in (v, w_0) on bounded subsets for any fixed $\omega \in \Omega$ and is measurable with respect to (v, w_0, ω) .

The proof of Lemma 1 is similar to that of Theorem 1 and is omitted. In what follows, we denote by $\mathcal{W}_t(v, w_0, \omega)$ the value of the function $\mathcal{W}(v, w_0, \omega)$ at time t .

To construct $\Psi_u : \Omega \rightarrow \Omega$, let us choose a smooth function θ on \mathbb{R} such that $\theta_t = 1$ for $t \leq 0$ and $\theta_t = 0$ for $t \geq 1$. For any $u \in H^1$, we set

$$F_N(u) = i\beta \mathbf{P}_N(|u|^{2\sigma}u), \quad v_t(\omega) = \mathbf{P}_N u_t(\omega, u), \quad w_t(\omega) = \mathbf{Q}_N u_t(\omega, u),$$

where $u_t(\omega, u)$ denotes the solution of (4), (5) issued from u . We now define Ψ_u by the relations

$$\mathbf{P}_N \Psi_u(\omega) = \mathbf{P}_N \omega - \mathbf{P}_N(\theta u - (\nu + i\alpha)\Theta \Delta u) + \int_0^t D_s(\omega, u) ds, \quad (40)$$

$$\mathbf{Q}_N \Psi_u(\omega) = \mathbf{Q}_N \omega, \quad (41)$$

where $\omega \in \Omega$, $\Theta_t = \int_0^t \theta_s ds$, and

$$D_s(\omega, u) = F_N(v_s - \theta_s \mathbf{P}_N u + \mathcal{W}_s(v - \theta \mathbf{P}_N u, 0, \omega)) - F_N(v_s + \mathcal{W}_s(v, \mathbf{Q}_N u, \omega)).$$

Step 2: Proof of (30). Lemma 1 implies that problem (4) – (6) is equivalent to the following system for the components (v_t, w_t) of the solution u_t :

$$\begin{aligned} \dot{v} - (\nu + i\alpha)\Delta v + F_N(v + \mathcal{W}_t(v, \mathbf{Q}u_0, \omega)) &= \mathbf{P}_N h + \mathbf{P}_N \eta, \\ \dot{w} - (\nu + i\alpha)\Delta w + F_N^\perp(v + w) &= \mathbf{Q}_N h + \mathbf{Q}_N \eta, \\ v_0 &= \mathbf{P}_N u_0, \quad w_0 = \mathbf{Q}_N u_0. \end{aligned}$$

Let us set

$$p_t = \mathbf{P}_N u_t(\Psi_u(\omega), 0), \quad q_t = \mathbf{Q}_N u_t(\Psi_u(\omega), 0).$$

Repeating literally the arguments used in [22] (see Step 7 in Section 3.1), we easily show that

$$p_t(\omega) = \mathbf{P}_N(u_t(\omega, u) - \theta_t u) \quad \text{for a.e. } \omega \in \Omega. \quad (42)$$

Combining this with (41), we see that, for almost every ω , the conditions of Proposition 5 with $T = 1$ (see Section 4.1) are satisfied for the pair of solutions $u_t(\omega, u)$ and $u_t(\Psi_u(\omega), 0)$. Therefore, combining (59) and (42), we obtain

$$\begin{aligned} \|u_t - \tilde{u}_t\|^2 &\leq C \exp\left(C \int_0^1 (\|u_s\|_1^{4\sigma} + \|\tilde{u}_s\|_1^{4\sigma}) ds\right) \|\mathbf{P}_N u\|^2 \\ &\quad + \exp\left(-\nu\alpha_{N+1}t + C \int_0^t (\|u_s\|_1^2 + \|\tilde{u}_s\|_1^2) ds\right) \|\mathbf{Q}_N u\|^2, \end{aligned} \quad (43)$$

where $\tilde{u}_t(\omega) = u_t(\Psi_u(\omega), 0)$ and $C > 0$ is a constant. Let us choose $N \geq 1$ so large that

$$\nu\alpha_{N+1} \geq 4C\nu^{-1}K, \quad (44)$$

where $K > 0$ is the constant in (15). Introduce the event

$$\Gamma(\rho, u) = \left\{ \int_0^1 (\|u_s\|_1^{4\sigma} + \|\tilde{u}_s\|_1^{4\sigma}) ds \geq \rho \text{ or } \sup_{t \geq 0} \int_0^t (\|u_s\|_1^2 + \|\tilde{u}_s\|_1^2) ds \geq \rho \right\}.$$

Propositions 1 and 2 and inequality (31) (which is proved below) imply that

$$\mathbb{P}(\Gamma(\rho, u)) \rightarrow 0 \quad \text{as } \rho \rightarrow \infty \text{ uniformly in } u \in B_{H^1}(1). \quad (45)$$

Furthermore, it follows from (15), (18) and (43) that, on the complement of $\Gamma(\rho, u)$, we have

$$\sup_{t \geq 0} \|u_t - \tilde{u}_t\| \leq C_\rho \|u\| \quad \text{for } u \in B_{H^1}(1),$$

where $C_\rho > 0$ does not depend on u . Combining this with (45), we see that inequality (30) holds with a constant $\delta = \delta(\varepsilon) > 0$ going to zero with ε .

Step 3: Proof of (31). We follow the scheme used in [4, 13]. Let us write Ψ_u in the form

$$\Psi_u(\omega) = \omega_t + \int_0^t \varphi_s(\omega, u) ds,$$

where φ_t is an H_N -valued function defined as

$$\varphi_t(\omega, u) = -\dot{\theta}_t \mathbf{P}_N u + (\nu + i\alpha)\theta_t \Delta \mathbf{P}_N u + D_t(\omega, u).$$

Introduce the functions

$$u_1 = v_t - \theta_t \mathbf{P}_N u + \mathcal{W}_t(v - \theta \mathbf{P}_N u, 0, \omega), \quad u_2 = v_t + \mathcal{W}_t(v, \mathbf{Q}_N u, \omega).$$

Let us fix a parameter $\rho > 0$ and define a truncating function χ^ρ by the following rule: $\chi_t^\rho(\omega, u) = 1$ if

$$\int_0^{t \wedge 1} (\|u_1\|_1^{4\sigma} + \|u_2\|_1^{4\sigma}) ds \leq \rho, \quad C(\mathcal{E}_{u_1}(t) + \mathcal{E}_{u_2}(t)) \leq 2Kt + \rho,$$

where K and C are the constants in inequality (15) and Proposition 5, respectively, and $\chi_t^\rho(\omega, u) = 0$ otherwise. Along with Ψ_u , let us consider the transformation

$$\Phi_u(\omega) = \omega_t + \int_0^t d_s(\omega, u) ds, \quad d_t(\omega, u) = \chi_t^\rho(\omega, u) \varphi_t(\omega, u).$$

In view of the triangle inequality and an elementary property of the total variation distance, we have

$$\begin{aligned} \|\mathbb{P} - \Psi_{u^*}(\mathbb{P})\|_{\text{var}} &\leq \|\mathbb{P} - \Phi_{u^*}^\rho(\mathbb{P})\|_{\text{var}} + \|\Phi_{u^*}^\rho(\mathbb{P}) - \Psi_{u^*}(\mathbb{P})\|_{\text{var}} \\ &\leq \|\mathbb{P} - \Phi_{u^*}^\rho(\mathbb{P})\|_{\text{var}} + \mathbb{P}\{\Phi_{u^*}^\rho \neq \Psi_{u^*}\}. \end{aligned} \quad (46)$$

Propositions 1 and 2 and the definition of χ^ρ imply that

$$\mathbb{P}\{\Phi_{u^*}^\rho \neq \Psi_{u^*}\} \rightarrow 0 \quad \text{as } \rho \rightarrow \infty \text{ uniformly in } u \in B_{H^1}(1). \quad (47)$$

To estimate the first term on the right-hand side of (46), we use Proposition 6 (see Section 4.1). We claim that if N is such that inequality (44) holds with a sufficiently large $C > 0$, then

$$\|d_t(\omega, u)\| \leq C_{N,\rho}(1+t)^r (\|u_1\| + \|u_2\|) (\theta_t \|\mathbb{P}_N u\| + e^{-ct} \|\mathbb{Q}_N u\|), \quad (48)$$

where $C_{N,\rho}$, c , and r are some positive constants not depending on u . If this inequality is established, then

$$\Lambda(N, \rho, u) := \exp\left(6 \int_0^\infty \|K^{-1} d_s(\omega, u)\|^2 ds\right) \leq \exp(C'_{N,\rho} \|u\|_1^2).$$

Thus, condition (63) is satisfied, and (64) implies that, for any fixed $\rho > 0$,

$$\|\mathbb{P} - \Phi_{u^*}^\rho(\mathbb{P})\|_{\text{var}} \leq \frac{1}{2} (\sqrt{\Lambda(N, \rho, u)} - 1)^{1/2} \rightarrow 0 \quad \text{as } \|u\|_1 \rightarrow 0.$$

Combining this with (46) and (47), we arrive at inequality (31) in which $\delta \rightarrow 0$ as $\|u\|_1 \rightarrow 0$. Thus, it remains to prove (48).

To this end, first note that

$$\|d_t(\omega, u)\| \leq C_1 \theta_t \|\mathbb{P}_N u\| + \chi_t^\rho(\omega, u) \|F_N(u_1) - F_N(u_2)\|, \quad (49)$$

where we denote by C_i , $i = 1, 2, \dots$, unessential positive constants that may depend on N and ρ . Since $\dim H_N < \infty$, we have

$$\begin{aligned} \|F_N(u_1) - F_N(u_2)\| &\leq C_2 \|F_N(u_1) - F_N(u_2)\|_{L^1} \\ &\leq C_3 \int_D (|u_1|^{2\sigma} + |u_2|^{2\sigma}) |u_1 - u_2| dx \\ &\leq C_3 \left(\|u_1\|_{L^{4\sigma}} + \|u_2\|_{L^{4\sigma}}^{2\sigma} \right) \|u_1 - u_2\|. \end{aligned} \quad (50)$$

Using the Sobolev embedding theorems and the Gagliardo–Nirenberg inequality, we easily show that if $\chi_t^\rho(\omega, u) = 1$, then

$$\|u_i\|_{L^{4\sigma}}^{2\sigma} \leq C_4 (1+t)^r \|u_i\|_1, \quad i = 1, 2,$$

where $r = r(\sigma, n) > 0$ is a constant. Combining this with (49) and (50), we derive

$$\|d_t(\omega, u)\| \leq C_5 \left(\theta_t \|\mathbb{P}_N u\| + (1+t)^r (\|u_2\|_1 + \|u_2\|_1) \|u_1 - u_2\| \right).$$

Substituting (59) into this inequality, we arrive at (48).

3.4. Proof of Proposition 4. Let us fix two measures $\mu_1, \mu_2 \in \mathcal{P}(H^1)$ that satisfy (34). In view of the Kantorovich–Rubinstein theorem (see Theorem 11.8.2 in [3]), there exist H^1 -valued random variables u_{i0} , $i = 1, 2$, with distribution μ_i such that

$$\mathbb{E} \|u_{10} - u_{20}\| \leq \delta. \quad (51)$$

Let us denote by $u_i(t, x)$, $i = 1, 2$, the solution of (4) – (6) with $u_0 = u_{i0}$.

Step 1. To prove (35), we first estimate $\|u_1(1) - u_2(1)\|_1$. Let us note that almost every realization of the function $u = u_1 - u_2$ belongs to the space \mathcal{X} and satisfies the deterministic equation

$$\dot{u} - (\nu + i\alpha)\Delta u + g(u_1, u_2) = 0, \quad (52)$$

where $g(u_1, u_2) = i\beta(|u_1|^{2\sigma}u_1 - |u_2|^{2\sigma}u_2)$. Moreover, the boundary and initial conditions (5) and (6) with $u_0 = u_{10} - u_{20}$ are also verified. Taking the scalar product in H of (52) with the function $2(1 - \nu t\Delta)u$, we obtain

$$\partial_t \mathcal{N}_0(u) - \nu \|u\|_1^2 + 2\nu \mathcal{N}_1(u) + 2(g(u_1, u_2), (1 - \nu t\Delta)u) = 0, \quad (53)$$

where we set

$$\mathcal{N}_0(u) = \|u\|^2 + \nu t \|u\|_1^2, \quad \mathcal{N}_1(u) = \|u\|_1^2 + \nu t \|u\|_2^2.$$

Using the Hölder inequality and the Sobolev embedding theorems, we derive

$$\begin{aligned} 2|(g(u_1, u_2), (1 - \nu t\Delta)u)| &\leq C_1 \int_D (|u_1|^{2\sigma} + |u_2|^{2\sigma}) (|u|^2 + \nu t |u| |\Delta u|) dx \\ &\leq C_2 (\|u_1\|_1^{2\sigma} + \|u_2\|_1^{2\sigma}) (\|u\| \|u\|_1 + \nu t \|u\|_1 \|u\|_2) \\ &\leq C_3 (\|u_1\|_1^{2\sigma} + \|u_2\|_1^{2\sigma}) \mathcal{N}_0(u) + \nu \mathcal{N}_1(u). \end{aligned}$$

Substitution of this inequality into (53) results in

$$\partial_t \mathcal{N}_0(u) \leq C_3 (\|u_1\|_1^{2\sigma} + \|u_2\|_1^{2\sigma}) \mathcal{N}_0(u).$$

Applying the Gronwall inequality, we obtain

$$\begin{aligned} \|u_1(1) - u_2(1)\|_1^2 &\leq \nu^{-1} \|u_{10} - u_{20}\|^2 \exp\left\{C_3 \int_0^1 (\|u_1\|_1^{2\sigma} + \|u_2\|_1^{2\sigma}) ds\right\} \\ &\leq \nu^{-1} D(u_1, u_2) \|u_{10} - u_{20}\|^2, \end{aligned} \quad (54)$$

where

$$D(u_1, u_2) = \exp\left\{C_3 \sup_{0 \leq t \leq 1} (\|u_1(t)\|_1^{2\sigma} + \|u_2(t)\|_1^{2\sigma})\right\}.$$

Step 2. We now take any function $f \in \mathcal{L}(H^1)$ with norm $\|f\|_{\mathcal{L}} \leq 1$ and note that

$$\begin{aligned} |(f, P_1^* \mu_1 - P_1^* \mu_2)| &= |\mathbb{E} \{f(u_1(1)) - f(u_2(1))\}| \\ &\leq \mathbb{E} (\|u_1(1) - u_2(1)\|_1 \wedge 2). \end{aligned} \quad (55)$$

Introduce the event

$$\Gamma_R = \{\mathcal{H}_1(u_{10}) \vee \mathcal{H}_1(u_{20}) \geq R \text{ or } D(u_1, u_2) \geq R\}.$$

It follows from the first inequality in (34) and Proposition 2 that $\mathbb{P}(\Gamma_R) \rightarrow 0$ as $R \rightarrow \infty$. Choosing a sufficiently large R , we conclude from (54) and (55) that

$$\begin{aligned} |(f, P_1^* \mu_1 - P_1^* \mu_2)| &\leq \frac{\delta}{2} + \mathbb{E} (I_{\Gamma_R} \|u_1(1) - u_2(1)\|_1) \\ &\leq \frac{\delta}{2} + \nu^{-1} R \mathbb{E} \|u_{10} - u_{20}\|. \end{aligned}$$

Combining this with (51) and recalling that f was arbitrary, we arrive at the required result.

4. Appendix.

4.1. Foias–Prodi type estimate. In this subsection, we study problem (4) – (6) in which η is a deterministic function of the form (cf. (7))

$$\eta(t, x) = \frac{\partial}{\partial t} \zeta(t, x), \quad (56)$$

where ζ is a continuous function from \mathbb{R}_+ to H^1 . Recall that $\{e_j\} \subset H$ is the complete set of eigenfunctions for the Dirichlet Laplacian in the domain D , H_N is the $2N$ -dimensional subspace generated by $\{e_j, ie_j, 1 \leq j \leq N\}$, and H_N^\perp is the orthogonal complement of H_N in H . Denote by $P_N : H \rightarrow H_N$ and $Q_N : H \rightarrow H_N^\perp$ the corresponding orthogonal projections.

The following result provides a Foias–Prodi type estimate for the difference between two solutions whose projections to H_N coincide (cf. [5]).

Proposition 5. *Let $n \geq 3$ and $\sigma \leq \frac{2}{n}$, let $h \in H^1(D, \mathbb{C})$, and let $u_1, u_2 \in \mathcal{X} = C(\mathbb{R}_+, H^1) \cap L_{\text{loc}}^2(\mathbb{R}_+, H^2)$ be two solutions of problem (4), (56) that correspond to deterministic functions $\zeta_1, \zeta_2 \in C(\mathbb{R}_+, H^1)$. Suppose that*

$$P_N u_1(t) = P_N u_2(t) \quad \text{for } t \geq T, \quad (57)$$

$$Q_N \zeta_1(t) = Q_N \zeta_2(t) \quad \text{for } t \geq 0, \quad (58)$$

where $T \geq 0$ and $N \geq 1$ is an integer. Then there is a constant $C > 0$ not depending on u_1, u_2, T , and N such that

$$\begin{aligned} \|Q_N(u_1(t) - u_2(t))\|^2 &\leq \exp\{-\nu\alpha_{N+1}t + q(t)\} \|Q_N(u_1(0) - u_2(0))\|^2 \\ &+ C e^{q(T)} \int_0^t (\|u_1(s)\|_1^{4\sigma} + \|u_2(s)\|_1^{4\sigma}) \|P_N(u_1(s) - u_2(s))\|_1^2 ds, \end{aligned} \quad (59)$$

where we set

$$q(t) = C \int_0^t (\|u_1(s)\|_1^2 + \|u_2(s)\|_1^2) ds.$$

Proof. The difference $u = u_1 - u_2$ belongs to the space \mathcal{X} and satisfies the equation

$$\dot{u} - (\nu + i\alpha)\Delta u + g(u_1, u_2) = \dot{\zeta}_1 - \dot{\zeta}_2, \quad (60)$$

where $g(u_1, u_2) = i\beta(|u_1|^{2\sigma}u_1 - |u_2|^{2\sigma}u_2)$. Let us set $v = P_N u$ and $w = Q_N u$ and take the scalar product in H of Eq. (60) with the function $2w$. Using (58), we derive

$$\partial_t \|w\|^2 + 2\nu \|w\|_1^2 + 2(g(u_1, u_2), w) = 0. \quad (61)$$

The Hölder inequality implies that

$$\begin{aligned} |(g(u_1, u_2), w)| &\leq C_1 \int_D (|u_1|^{2\sigma} + |u_2|^{2\sigma}) (|w|^2 + |v||w|) dx \\ &\leq C_1 \left(\|u_1\|_{L^{2q\sigma}}^{2\sigma} + \|u_2\|_{L^{2q\sigma}}^{2\sigma} \right) \left(\|w\|_{L^{2p}}^2 + \|v\|_{L^r} \|w\|_{L^{\frac{2n}{n-2}}} \right), \end{aligned} \quad (62)$$

where we set

$$p = \frac{n}{n(1-\sigma) + 2\sigma}, \quad q = \frac{n}{\sigma(n-2)}, \quad r = \frac{2n}{n+2-2\sigma(n-2)}.$$

In view of the Sobolev embedding theorems and the Gagliardo–Nirenberg inequality (see Theorem 6.4.1 in [7]), we have

$$\|w\|_{L^{2p}} \leq C_2 \|w\|^\theta \|w\|_1^{1-\theta}, \quad \|v\|_{L^r} \leq C_3 \|v\|_1,$$

where $\theta = \sigma + 1 - \frac{n\sigma}{2}$. Substitution of these inequalities into (62) results in

$$\begin{aligned} |(g(u_1, u_2), w)| &\leq C_4(\|u_1\|_1^{2\sigma} + \|u_2\|_1^{2\sigma}) (\|w\|^{2\theta} \|w\|_1^{2(1-\theta)} + \|v\|_1 \|w\|_1) \\ &\leq \frac{\nu}{2} \|w\|_1^2 + C_5(\|u_1\|_1^2 + \|u_2\|_1^2) \|w\|^2 + C_6(\|u_1\|_1^{4\sigma} + \|u_2\|_1^{4\sigma}) \|v\|_1^2. \end{aligned}$$

Combining this with (61) and using the inequality $\|w\|_1^2 \geq \alpha_{N+1} \|w\|^2$, we obtain

$$\partial_t \|w\|^2 + (\nu\alpha_{N+1} - 2C_5(\|u_1\|_1^2 + \|u_2\|_1^2)) \|w\|_1^2 \leq 2C_6(\|u_1\|_1^{4\sigma} + \|u_2\|_1^{4\sigma}) \|v\|_1^2.$$

Applying the Gronwall inequality and recalling that $v(t) = 0$ for $t \geq T$ in view of (57), we arrive at (59). \square

4.2. Girsanov theorem. Let $(\Omega, \mathcal{F}, \mathbb{P})$ be the probability space described in the beginning of Section 3.3 and let $\{\mathcal{F}_t, t \geq 0\}$ be the natural filtration on Ω augmented with respect to $(\mathcal{F}, \mathbb{P})$ (see [8]). Suppose that $N \geq 1$ is an integer and $d_t(\omega)$ is an H_N -valued measurable function on $\mathbb{R}_+ \times \Omega$ that is adapted to \mathcal{F}_t and satisfies the condition

$$\mathbb{E} \int_0^t \|d_s\|^2 ds < \infty \quad \text{for any } t \geq 0.$$

Consider a transformation of Ω defined as

$$\Phi(\omega) = \omega + \int_0^t d_s(\omega) ds.$$

Let K be a bounded linear operator in H defined by the relations $Ke_j = b_j e_j$ and $K(ie_j) = ib_j e_j$ for all $j \geq 1$. In view of condition (28), the operator K is injective and its inverse K^{-1} is well defined on H_N . The following result is a straightforward consequence of the Girsanov theorem (see [20]).

Proposition 6. *Suppose that*

$$\Lambda := \mathbb{E} \exp\left(6 \int_0^\infty \|K^{-1}d_s\|^2 ds\right) < \infty. \quad (63)$$

Then the total variation distance between \mathbb{P} and $\Phi_(\mathbb{P})$ admits the estimate*

$$\|\mathbb{P} - \Phi_*(\mathbb{P})\|_{\text{var}} \leq \frac{1}{2} (\sqrt{\Lambda} - 1)^{1/2}. \quad (64)$$

Proof. We define a Brownian motion in H_N by the formula $B_t(\omega) = K^{-1}\mathbb{P}_N\omega_t$ and introduce the measurable function

$$\rho(\omega) = \exp\left(-\int_0^\infty K^{-1}d_s(\omega) dB_s - \frac{1}{2} \int_0^\infty \|K^{-1}d_s(\omega)\|^2 ds\right).$$

In view of the Girsanov theorem (see Theorem 8.6.4 in [20]), the function ρ is the density (with respect to \mathbb{P}) of a probability measure on Ω , and the random process

$$\widehat{B}_t(\omega) = B_t(\omega) + \int_0^t K^{-1}d_s(\omega) ds$$

is a Brownian motion with respect to the measure $\widehat{\mathbb{P}}(d\omega) = \rho(\omega)\mathbb{P}(d\omega)$. It follows that the distribution of $\Phi(\omega)$ under the law $\widehat{\mathbb{P}}$ coincides with \mathbb{P} . Therefore,

$$\begin{aligned} \|\mathbb{P} - \Phi_*(\mathbb{P})\|_{\text{var}} &= \sup_{\Gamma \in \mathcal{F}} |\mathbb{P}(\Gamma) - \mathbb{P}\{\Phi(\cdot) \in \Gamma\}| \\ &= \sup_{\Gamma \in \mathcal{F}} |\widehat{\mathbb{P}}\{\Phi(\cdot) \in \Gamma\} - \mathbb{P}\{\Phi(\cdot) \in \Gamma\}| \leq \|\widehat{\mathbb{P}} - \mathbb{P}\|_{\text{var}}. \end{aligned}$$

Using a well-known formula for the total variation distance, we derive

$$\|\widehat{\mathbb{P}} - \mathbb{P}\|_{\text{var}} = \frac{1}{2} \int_{\Omega} |\rho - 1| d\mathbb{P} \leq \frac{1}{2} \left(\int_{\Omega} (\rho - 1)^2 d\mathbb{P} \right)^{\frac{1}{2}} \leq \frac{1}{2} \left(\int_{\Omega} \rho^2 d\mathbb{P} - 1 \right)^{\frac{1}{2}}. \quad (65)$$

It remains to note that the integral on the right-hand side of (65) can be estimated by $\sqrt{\Lambda}$. \square

REFERENCES

- [1] J. Bricmont, *Ergodicity and mixing for stochastic partial differential equations*, Proceedings of the International Congress of Mathematicians, vol. 1, Higher Ed. Press, Beijing, 2002.
- [2] A. Debussche and C. Odasso, *Ergodicity for the weakly damped stochastic nonlinear Schrödinger equation*, J. Evol. Equ. (2005), to appear.
- [3] R. M. Dudley, *Real Analysis and Probability*, Cambridge University Press, Cambridge, 2002.
- [4] W. E, J. C. Mattingly, and Ya. G. Sinai, *Gibbsian dynamics and ergodicity for the stochastically forced Navier–Stokes equation*, Comm. Math. Phys. **224** (2001), 83–106.
- [5] C. Foias and G. Prodi, *Sur le comportement global des solutions non-stationnaires des équations de Navier–Stokes en dimension 2*, Rend. Sem. Mat. Univ. Padova **39** (1967), 1–34.
- [6] M. Hairer, *Exponential mixing properties of stochastic PDE’s through asymptotic coupling*, Probab. Theory Related Fields **124** (2002), 345–380.
- [7] L. Hörmander, *Lectures on Nonlinear Hyperbolic Differential Equations*, Springer-Verlag, Berlin, 1997.
- [8] I. Karatzas and S. E. Shreve, *Brownian Motion and Stochastic Calculus*, Springer-Verlag, New York, 1991.
- [9] N. V. Krylov, *SPDEs in $L_q((0, \tau], L_p)$ spaces*, Electron. J. Probab. **5** (2000), no. 13, 29 pp.
- [10] S. B. Kuksin, *Ergodic theorems for 2D statistical hydrodynamics*, Rev. Math. Phys. **14** (2002), 585–600.
- [11] S. Kuksin and A. Shirikyan, *Stochastic dissipative PDE’s and Gibbs measures*, Comm. Math. Phys. **213** (2000), no. 2, 291–330.
- [12] S. Kuksin and A. Shirikyan, *A coupling approach to randomly forced nonlinear PDE’s. I*, Comm. Math. Phys. **221** (2001), 351–366.
- [13] S. Kuksin and A. Shirikyan, *Coupling approach to white-forced nonlinear PDE’s*, J. Math. Pures Appl. **81** (2002), 567–602.
- [14] S. Kuksin and A. Shirikyan, *Randomly forced CGL equation: stationary measures and the inviscid limit*, J. Phys. A **37** (2004), no. 12, 3805–3822.
- [15] N. Masmoudi and L.-S. Young, *Ergodic theory of infinite dimensional systems with applications to dissipative parabolic PDE’s*, Comm. Math. Phys. (2002), 461–481.
- [16] J. Mattingly, *Exponential convergence for the stochastically forced Navier–Stokes equations and other partially dissipative dynamics*, Comm. Math. Phys. **230** (2002), 421–462.
- [17] P.-A. Meyer, *Probability and Potentials*, Balisdell, Waltham-Toronto-London, 1966.
- [18] R. Mikulevicius and B. Rozovskii, *A note on Krylov’s L_p -theory for systems of SPDEs*, Electron. J. Probab. **6** (2001), no. 12, 35 pp.
- [19] C. Odasso, *Ergodicity for the stochastic complex Ginzburg–Landau equations*, Ann. Inst. H. Poincaré Probab. Statist. (2005), to appear.
- [20] B. Øksendal, *Stochastic Differential Equations*, Springer-Verlag, Berlin, 1998.
- [21] A. Shirikyan, *Some mathematical problems of statistical hydrodynamics*, Proceedings of the XIV International Congress of Mathematical Physics, World Scientific, 2003, to appear.
- [22] A. Shirikyan, *Ergodicity for a class of Markov processes and applications to randomly forced PDE’s. I*, Russian J. Math. Phys. **12** (2005), no. 1, 81–96.

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E-mail address: Armen.Shirikyan@math.u-psud.fr